



# The topological insulator with strong correlation effects

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Lu et al, PRL 110, 096401 (2013)



# outline

- Band inversion and TI
- Mix valence compound: Band inversion between d and f bands; the important role of e-e interaction
- SmB<sub>6</sub> , YbB<sub>6</sub>, and YbBI<sub>2</sub>
- Conclusion

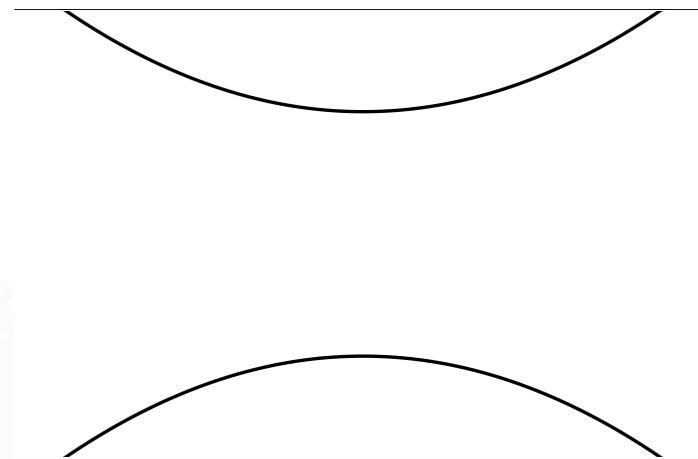
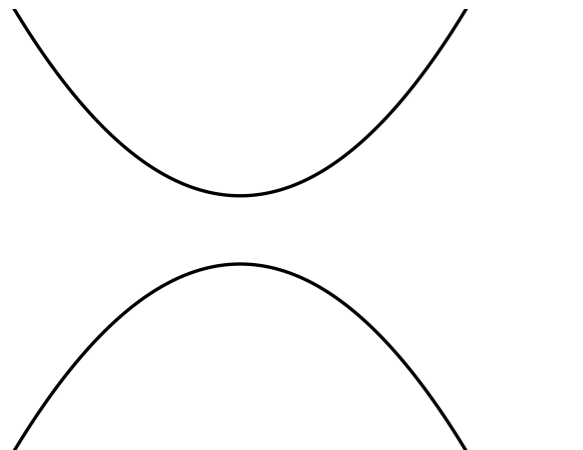


## Topological classification of band insulators

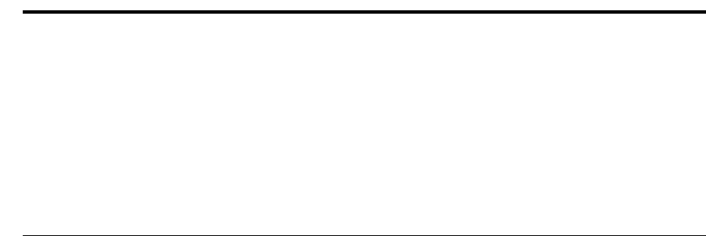
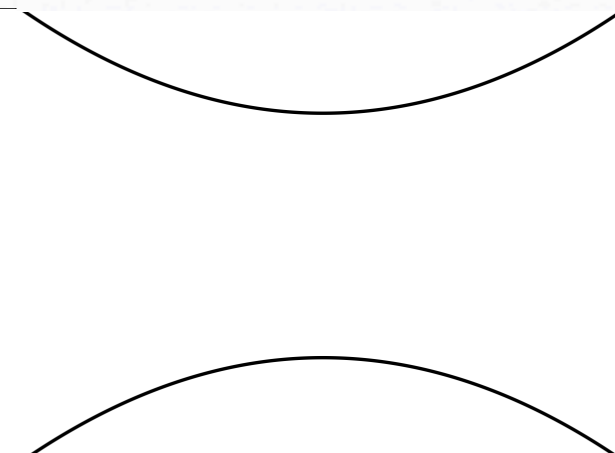
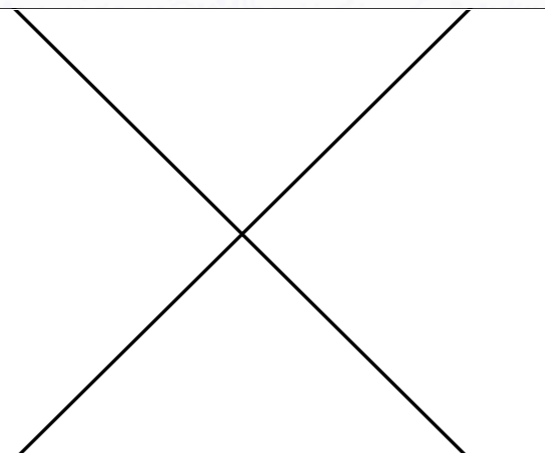
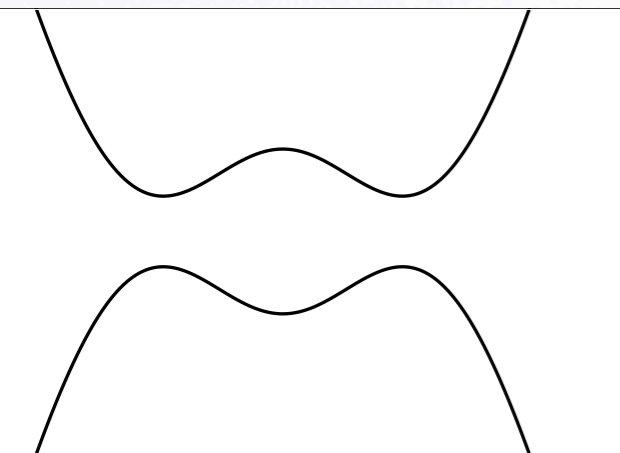
- normal insulator: the electronic structure can be smoothly transformed to isolated atoms
- can not be smoothly transformed to isolated atoms without going through a phase transition,  $Z_2$  invariance defined for non-interacting TI
- TI with interaction can be defined by the Theta angle of of topological magneto-electric effect



# Evolution of band structure from solid to isolated atoms



**Normal Insulator**



**Topological Insulator**

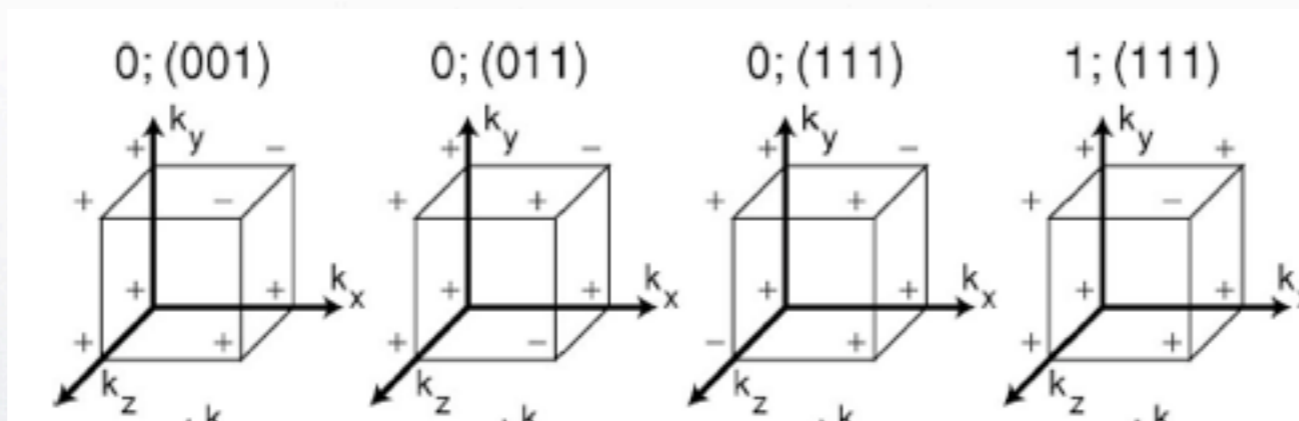


# Z2 invariance for non-interacting TI with inversion symmetry

- **Simple rules for TI with inversion symmetry:  
strong index and weak indices;  $\mathbf{K}=-\mathbf{K}$  high symmetry points**

$$(-1)^{\nu_0} = \prod_{i=1}^8 \delta_i.$$

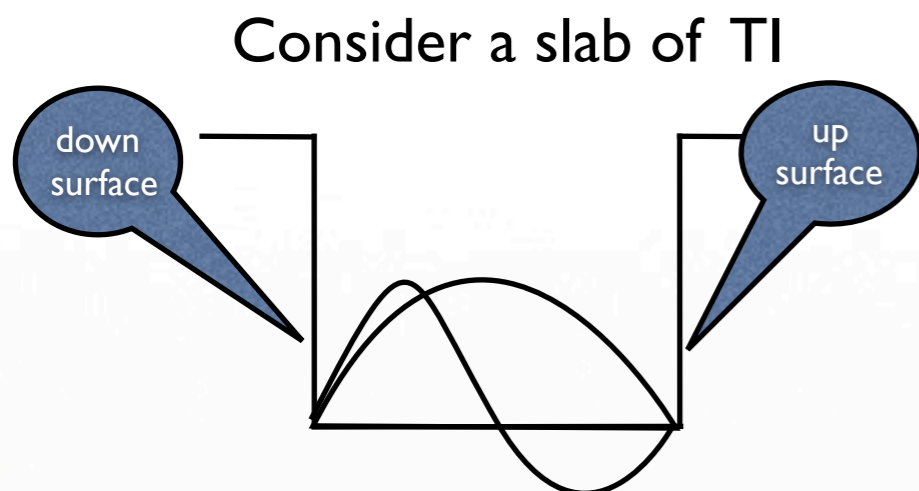
$$\delta_i = \prod_{m=1}^N \xi_{2m}(\Gamma_i).$$



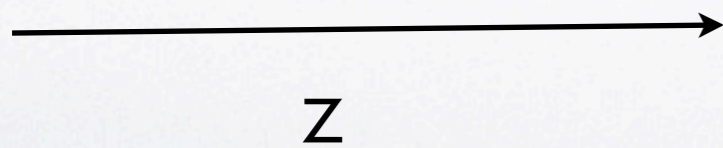
L. Fu and C. Kane, PRB 76,045302



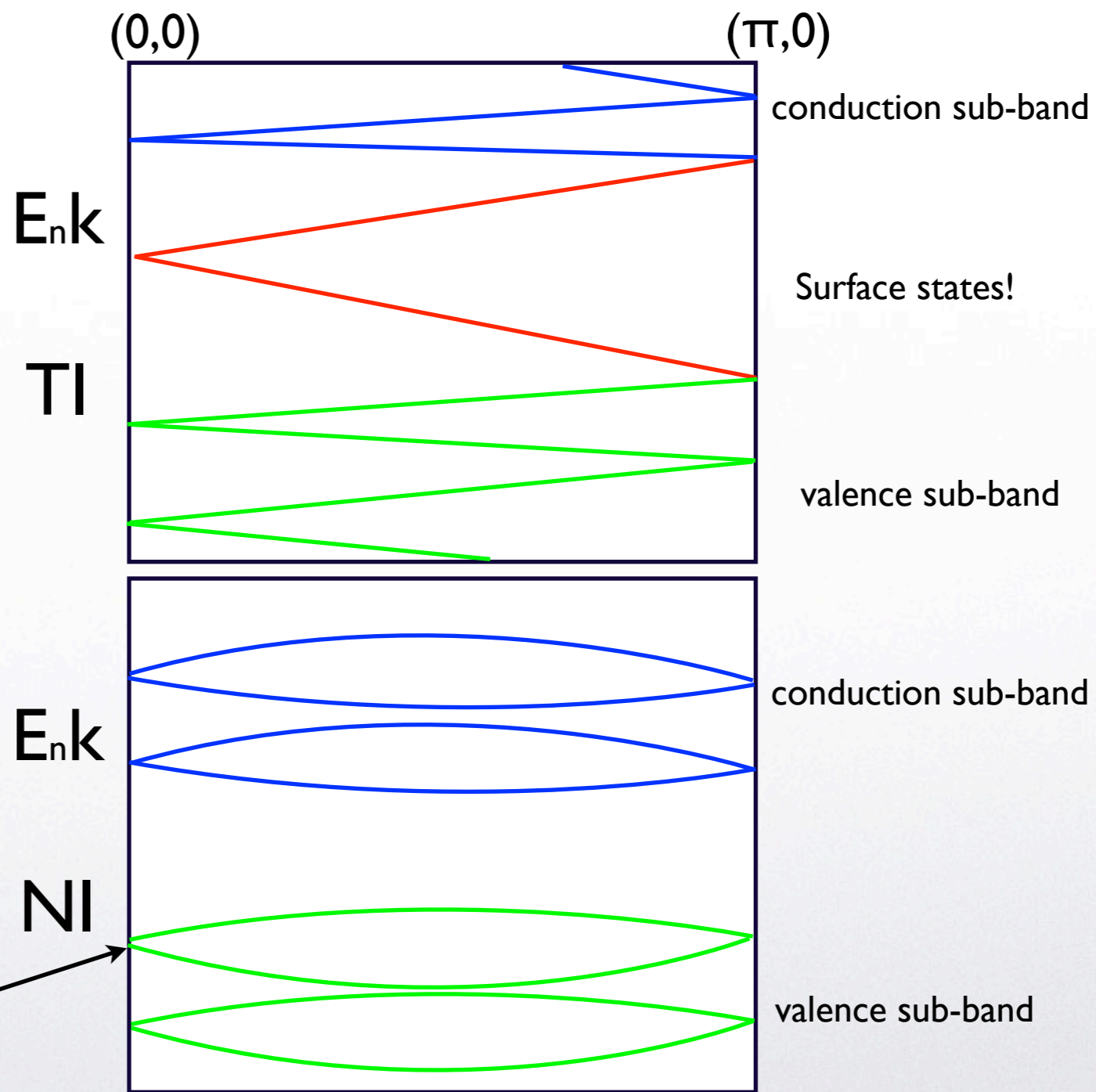
# TI and surface states



$$\psi_n(k_x, k_y)$$

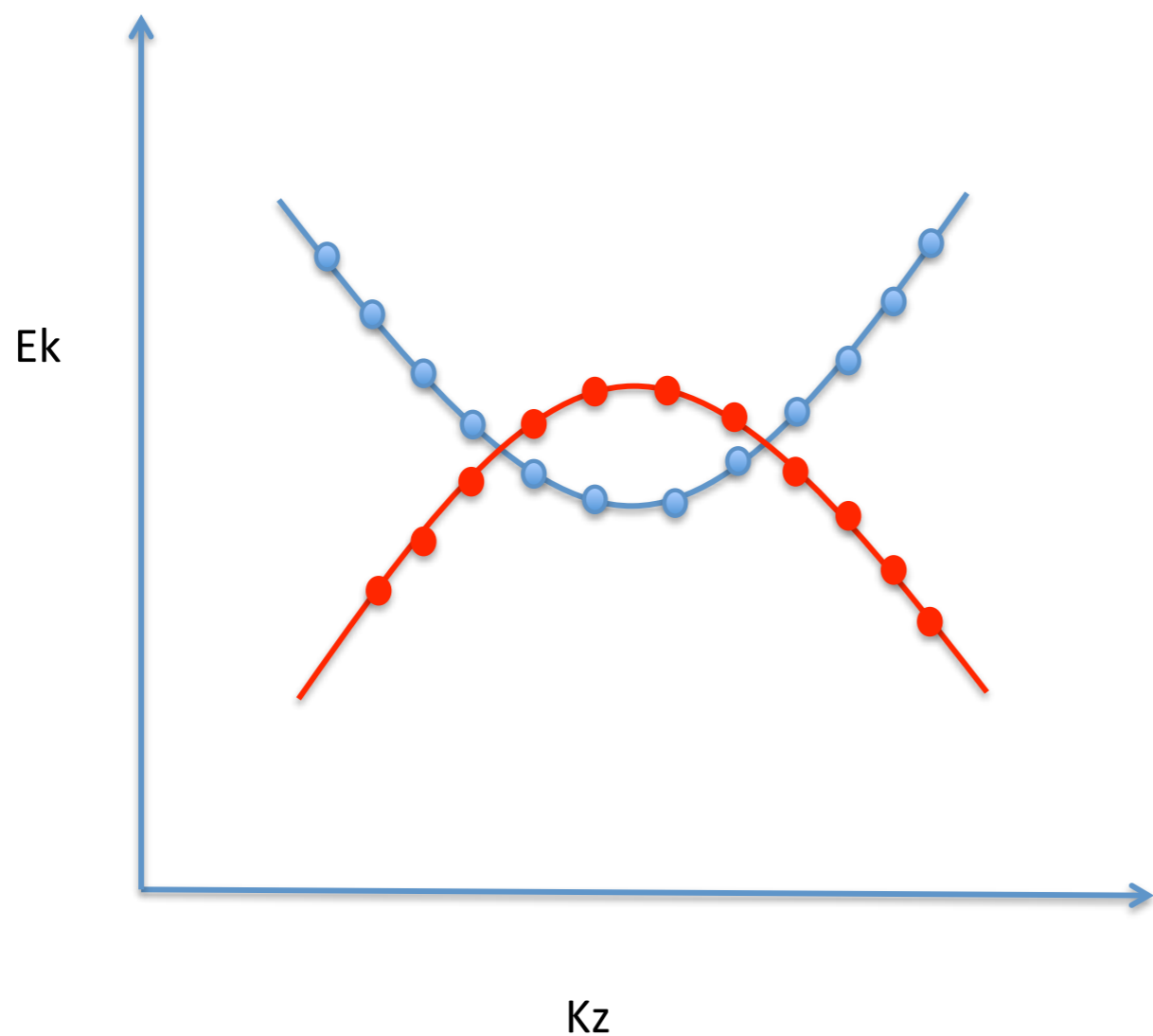


$k = -k + G$ ; Kramers degeneracy





# Introduction: Band Inversion and TI





# The different types of band inversion

- Band inversion between s and p bands: HgTe
- Band inversion between bonding and anti-bonding p-bands: Bi<sub>2</sub>Se<sub>3</sub>, Sb
- Band inversion generated by valence fluctuation of 4f/5f bands
- Very strong correlation effects in f-electron materials:
- Does these materials topologically non-trivial?



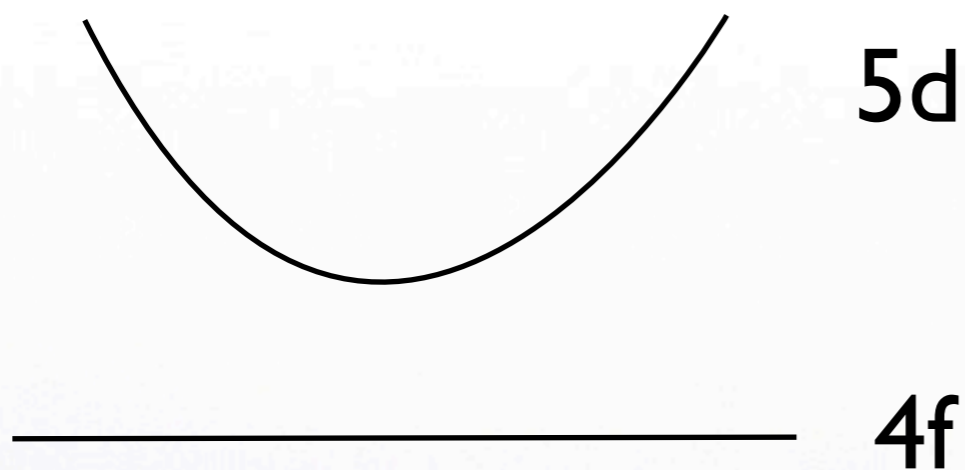


## 4f/5f compounds with intermediate valence

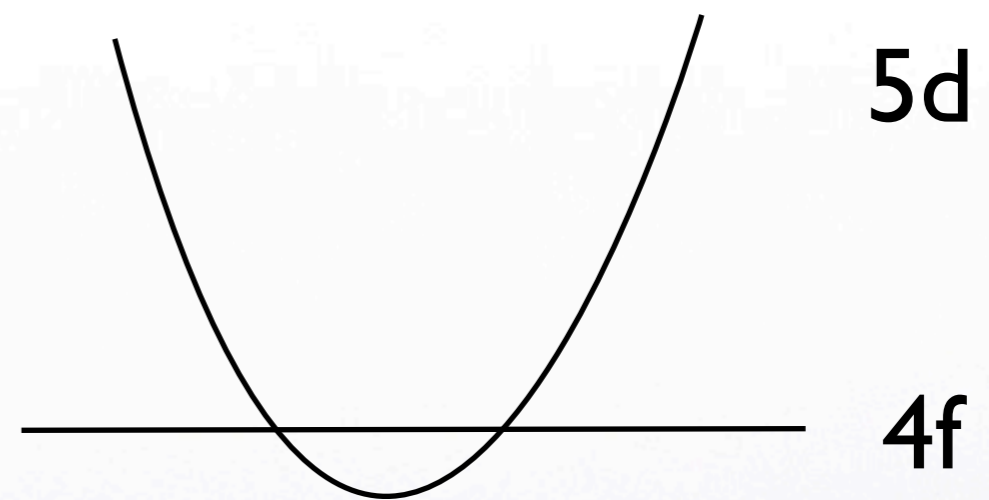
- golden phase of SmS
- SmB<sub>6</sub>, YbB<sub>6</sub>, YbBI<sub>2</sub>
- PuTe and PuSe
- Are these materials topological Kondo insulator? M. Dzero et al, Phys. Rev. Lett. 104, 119901, 2010



# Intermediate Valence and band inversion: from the band theory point of view



Rare earth compounds  
with divalence (**not very stable**),  
Sm, Eu, Yb.....



intermediate valence state  
band theory point of view



# Is the band description correct for these compounds?

- from non-interacting band insulator to strong coupling “Kondo insulator”
- how to capture the correct electronic structure?
- how to describe its topological nature?



# How to compute the Z2 invariance for interacting system?

Formula derived from the topological field theory:  
[PRL.105,256803 \(2010\)](#)

$$P_3 = \frac{\pi}{6} \int_0^1 du \int \frac{d^4k}{(2\pi)^4} \text{Tr} \epsilon^{\mu\nu\rho\sigma} [G\partial_\mu G^{-1} G\partial_\nu G^{-1} \times G\partial_\rho G^{-1} G\partial_\sigma G^{-1} G\partial_\mu G^{-1}] \quad (1)$$

- Pole expansion of the self energy without k dependence [PRB 85, 235135 \(2012\)](#); [EPL98 \(2012\) 57001](#)
- Using the eigenstate of  $H_0 + \Sigma(0)$ , condition: there is no singularity along the imaginary axis of self energy, [PRB 85, 165126 \(2012\)](#)

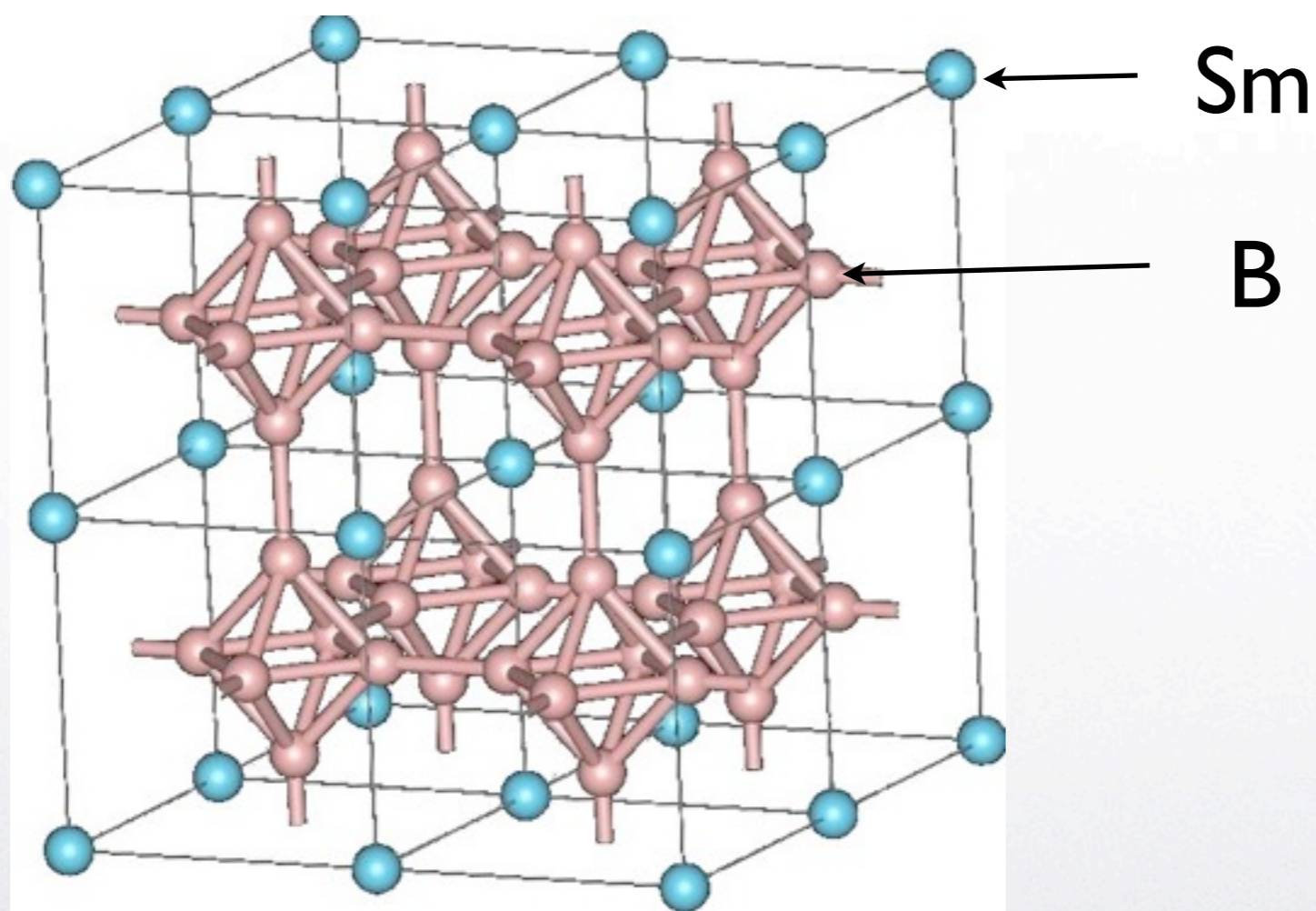


## Our method: LDA+Gutzwiller

- Gutzwiller type of ground state wave function
- Combined with first principle code
- Suitable for the study of f-electrons
- Has been applied to many correlated system:  
LaOFeAs;  $\text{NaxCoO2}$ ; Ce; Pu....PHYSICAL REVIEW B 79, 075114 ,2009



# The structure of SmB<sub>6</sub>





from PHYSICAL REVIEW B 66, 165209, 2002

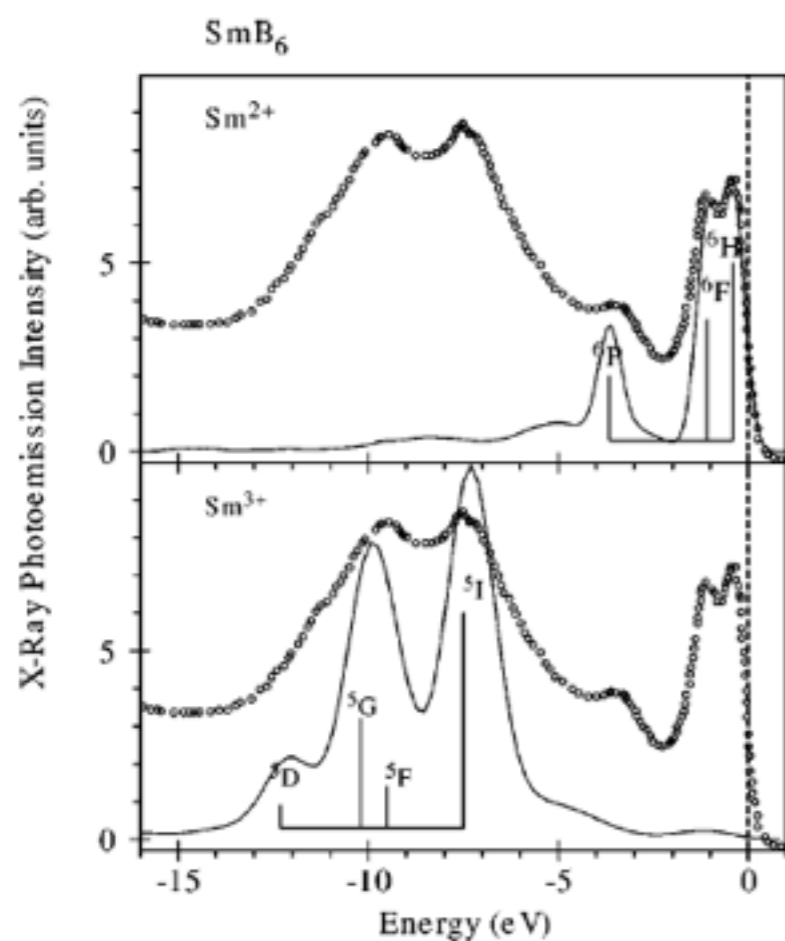
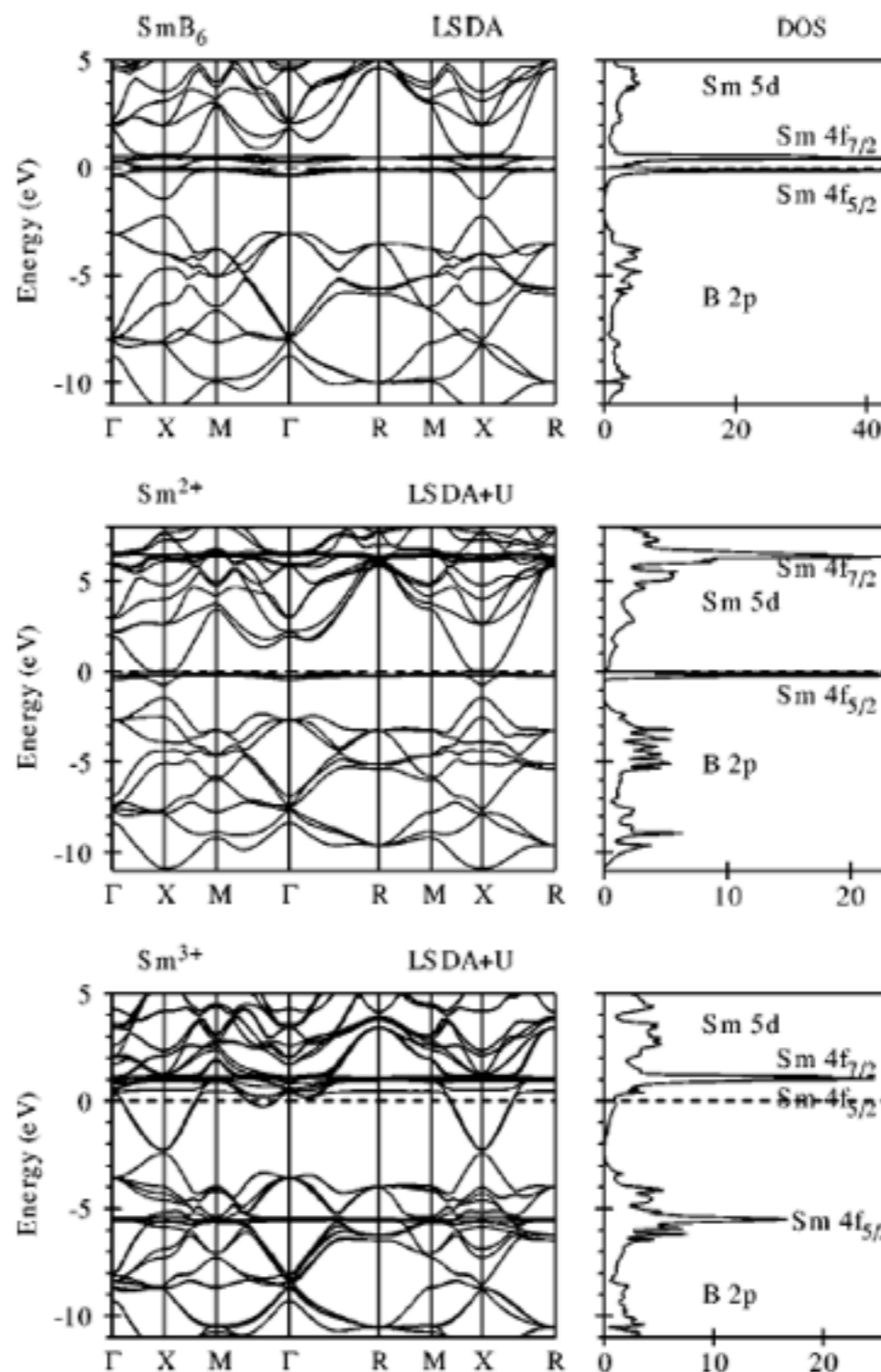


FIG. 3. Comparison of the calculated 4f DOS of SmB<sub>6</sub> using the LSDA+U approximation with the experimental XPS spectra from Ref. 18, taking into account the multiplet structure of the 4f<sup>5</sup> and 4f<sup>4</sup> final states (see explanations in the text).

compare to XPS

The valence determined by XPS is 2.54



band structure  
obtained by LDA  
and LDA+U

Assuming some  
kind of orbital  
ordering from Sm<sup>3+</sup>  
phase



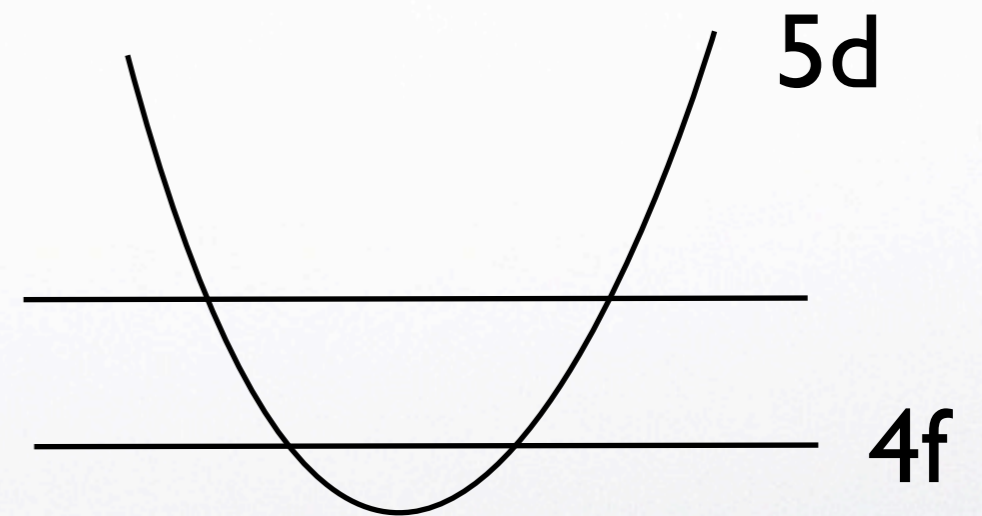
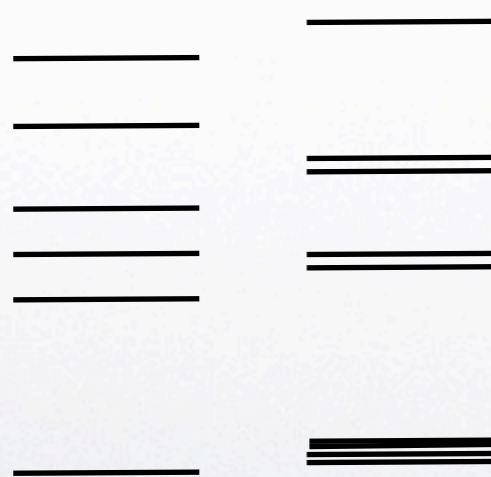
# Atomic multiplets or Band states?

$$H_{total} = H_{atom} + H_{hopping}$$

$$H_{atom} = F_0 \frac{(n_f - 1)n_f}{2} + \sum_{i=1, \alpha\beta\gamma\delta}^3 F_{2i}^{\alpha\beta\gamma\delta} f_{\alpha}^{\dagger} f_{\beta}^{\dagger} f_{\delta} f_{\gamma} + \sum_{\alpha} \epsilon_{\alpha} f_{\alpha}^{\dagger} f_{\alpha}$$

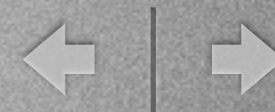
$$H_{hopping} = \sum_{ij\alpha\beta} t_{ij}^{\alpha\beta} f_{i\alpha}^{\dagger} f_{j\beta} + h.c.$$

... nf=6    nf=5    ...

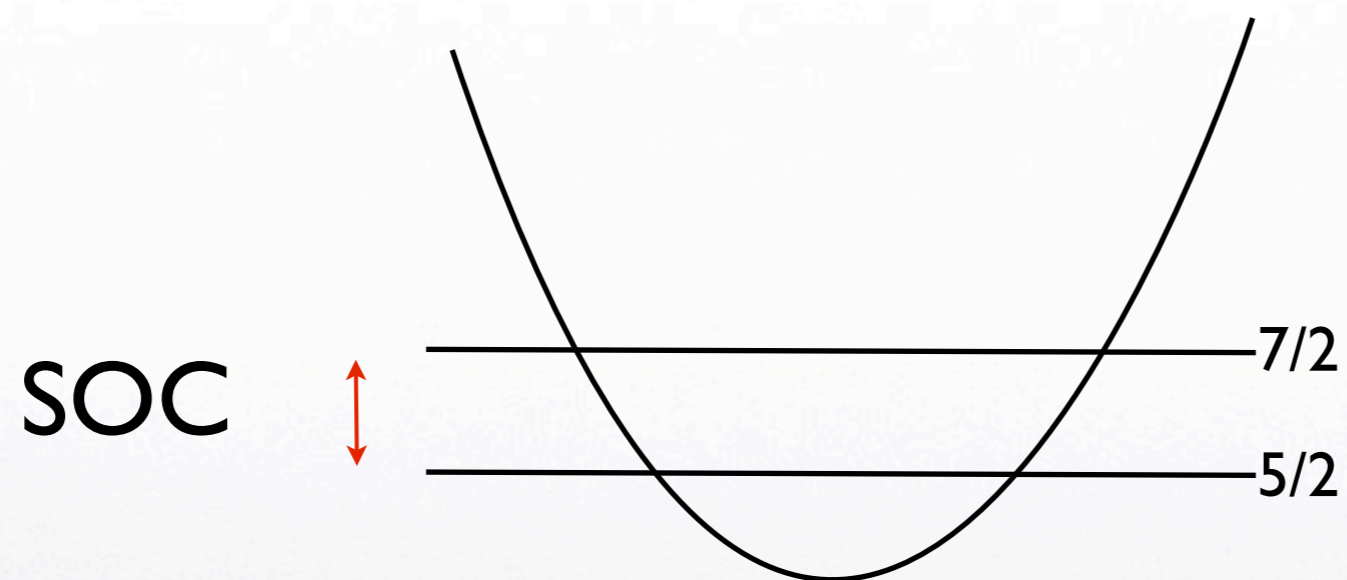


Both are non-fock states!



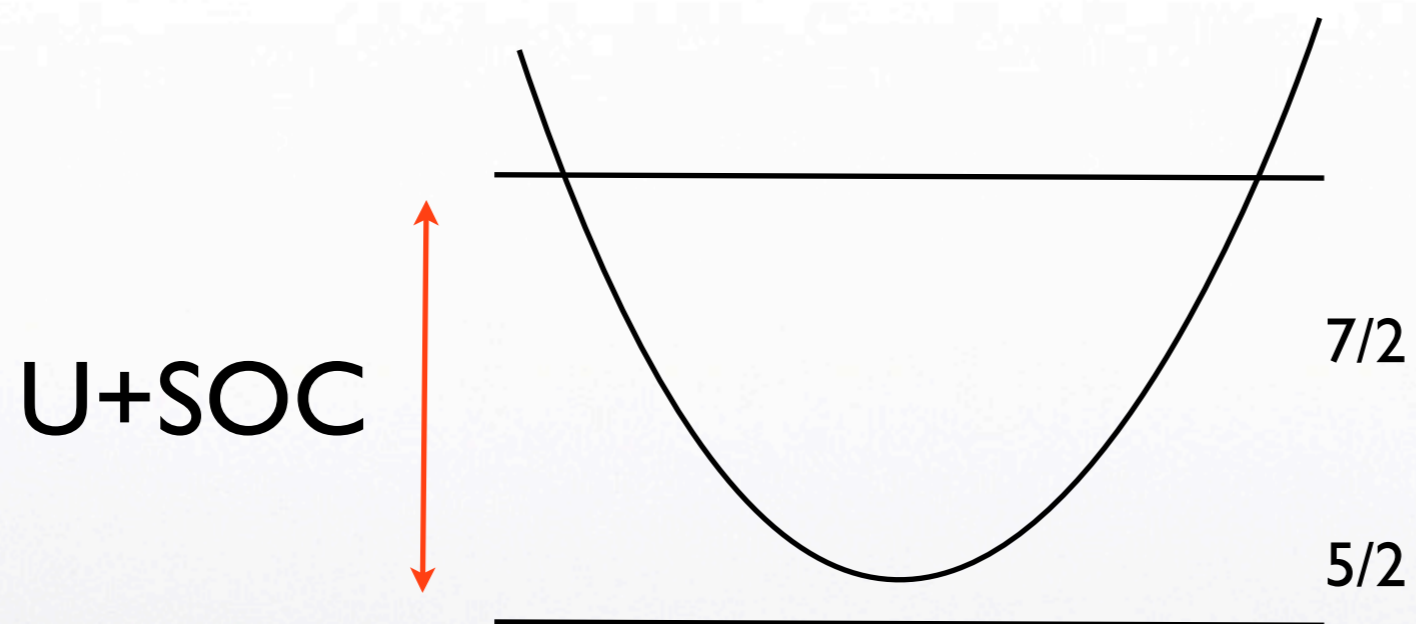


# Problem of Hartree-Fock treatment

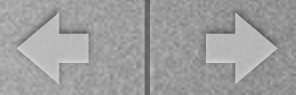




# Problem of Hartree-Fock treatment



MF atomic state: 



# The Gutzwiller Trial wave function used in this study

## Rotational Invariant Gutzwiller Approximation

Gutzwiller variational wavefunction:

$$|\Psi_G\rangle = \mathcal{P}|\Psi_0\rangle = \prod_{\mathbf{R}} \mathcal{P}_{\mathbf{R}}|\Psi_0\rangle$$

$$\mathcal{P}_{\mathbf{R}} = \sum_{\Gamma'} \lambda(\mathbf{R})_{\Gamma\Gamma'} |\Gamma, \mathbf{R}\rangle \langle \Gamma', \mathbf{R}|$$

$|\Gamma\rangle$ : eigenstates of atomic hamiltonian  $H_U$

$\Psi_0$ : uncorrelated wave function (Wick's Theorem holds)

$\mathcal{P}_{\mathbf{R}}$ : projector operator modify weight of local configuration

Gutzwiller Constraints:

$$\langle \Psi_0 | \mathcal{P}^\dagger \mathcal{P} | \Psi_0 \rangle = 1$$

$$\langle \Psi_0 | \mathcal{P}^\dagger \mathcal{P} n_{i\alpha} | \Psi_0 \rangle = \langle \Psi_0 | n_{i\alpha} | \Psi_0 \rangle$$

We only apply truncation respect to the occupation number, for SmB6, we keep all the atomic states with  $n_f=5,6,7$ , about 8000 variational parameters!



## Effective Hamiltonian in Gutzwiller approximation

- Under Gutzwiller approximation, we can define

$$E_G = \langle 0 | P H_{LDA} P | 0 \rangle + E_{int} \approx \langle 0 | H_{eff} | 0 \rangle + \sum_{\Gamma} \lambda_{\Gamma, \Gamma} E_{\Gamma}$$

It can be easily proved that  $H_{eff}$  is equivalent to  $H_0 + \Sigma(0)$  by comparing the Green's function in low frequency limit

$$G(i\omega) = \frac{z}{i\omega - H_{eff}} = \frac{1}{i\omega/z - H_{eff}/z}$$

$$H_0 + \Sigma(0) = -G^{-1}(0) = H_{eff}/z$$



## Main difficulties for the electronic structure calculation for f-electron systems

- Strong interactions among f-electrons, which can be expressed in terms of Slater integrals:  $F_0, F_2, F_4, F_6$
- for SmB<sub>6</sub>  $F_0=5.8\text{eV}$ ,  $F_2=9.9\text{eV}$ ,  $F_4=7.09\text{eV}$ ,  $F_6=4.99\text{eV}$
- Typical interaction strength is one order bigger than the f-band width
- Multiplet state VS the band state



## The double counting problem

- The total Hamiltonian treated in LDA+Gutzwiller

$$H_{total} = H_{LDA} + H_U + H_{DC}$$

$$H_{DC} = V_{DC} \sum_{k\sigma} f_{k\sigma}^\dagger f_{k\sigma}$$

- A commonly used form for DC term:

$$V_{ab}^{DC} = \delta_{ab} \left[ \bar{U} \left( \bar{n}_c - \frac{1}{2} \right) - \bar{J} \left( \bar{n}_c^\sigma - \frac{1}{2} \right) \right].$$

- Kotliar et al, RMP78, 865,2006



# Valence of Sm determined by exp

	paper	experiment	nf	Average valence
SmB6	PRB: 1976 14, 4586	XPS	5.3	2.7
SmB6	JAP: 1970 41, 898	Mossbauer effect	5.4	2.6
SmB6	Physica B: 1995 215, 99	neutron experiments	5.44	2.56
SmB6	JPCS: 2009 176, 012034	XAS	5.47	2.53

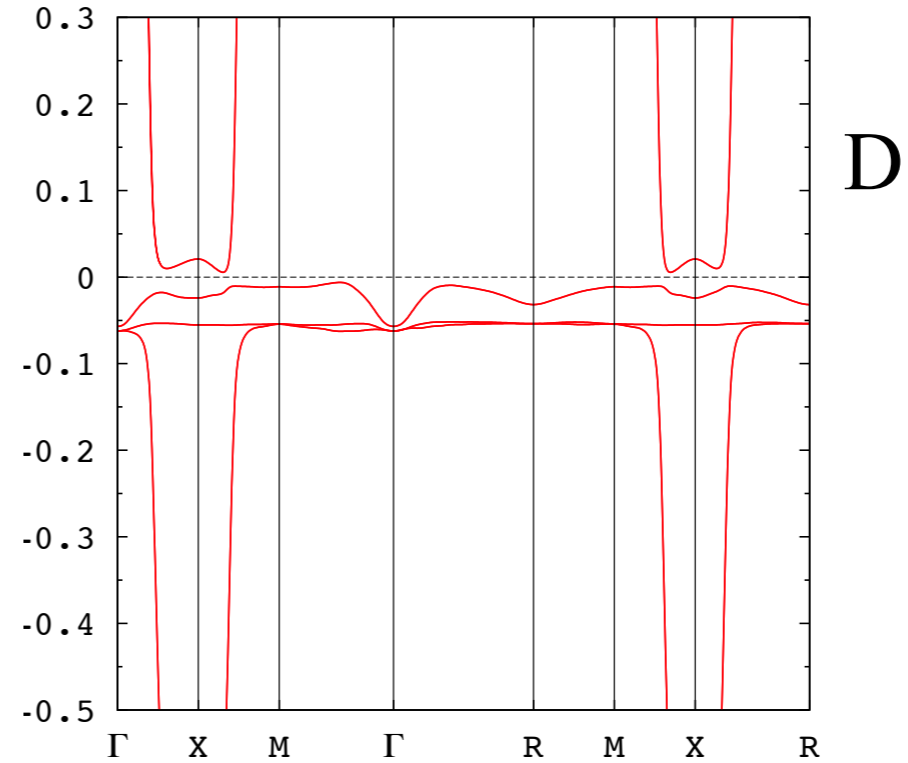
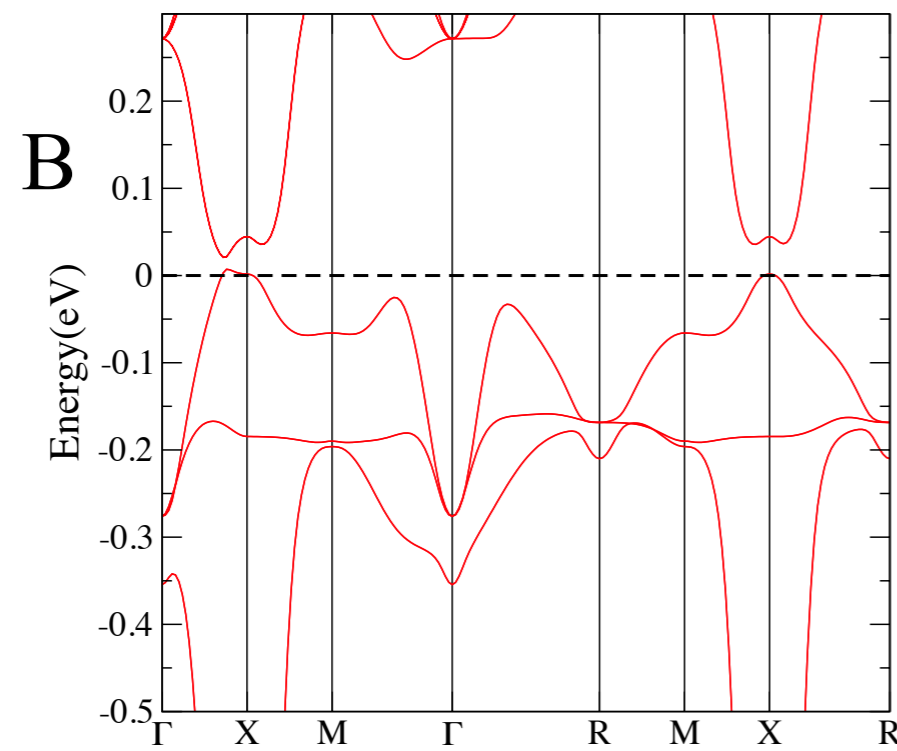
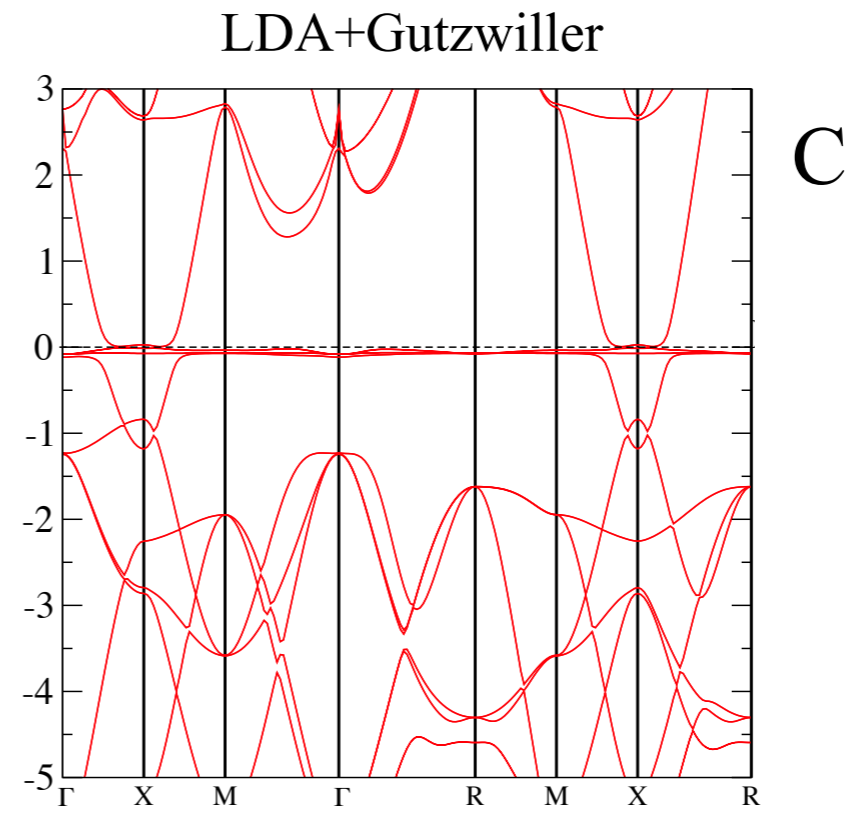
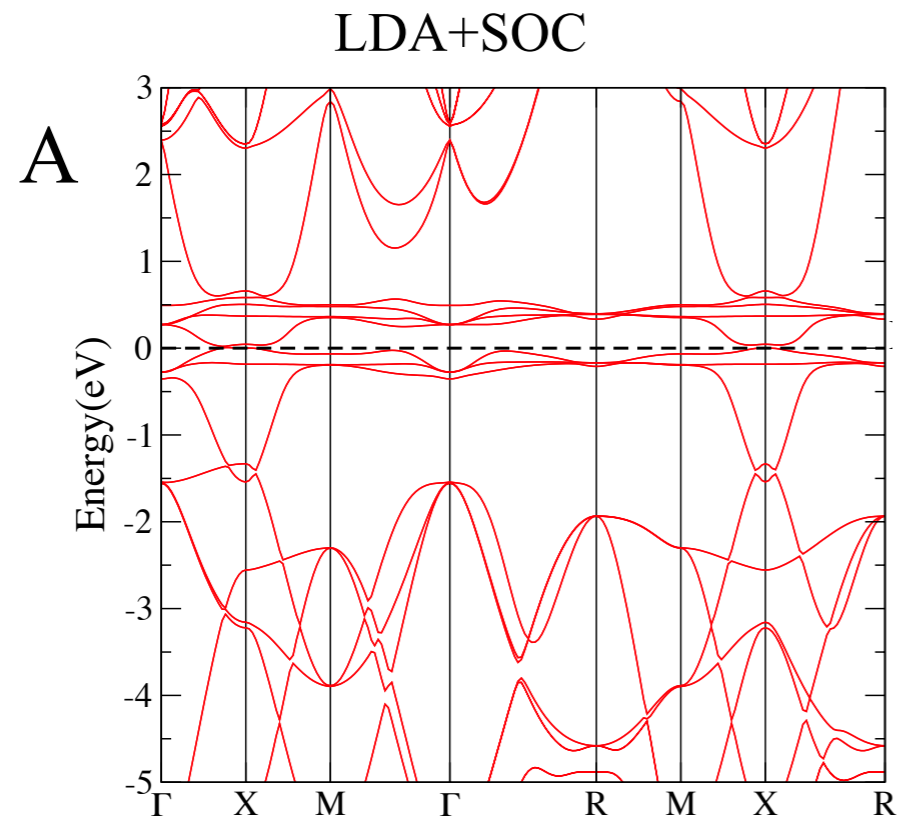
# SmB6: Z, nf, Gap vs Parity

		$E_g$	$\Gamma$	3X	R	3M	Tol
SmB6	LDA	25mev	+	-	+	+	-
<b>SmB6</b>	<b>LDA + G</b>	<b>10mev</b>	<b>+</b>	<b>-</b>	<b>+</b>	<b>+</b>	<b>-</b>

Vdc	$n_f$	$Z(R^2_{mat})$	
		5/2	7/2
	5.35		
<b>26.4</b>	<b>5.45</b>	<b>0.18</b>	<b>0.59</b>

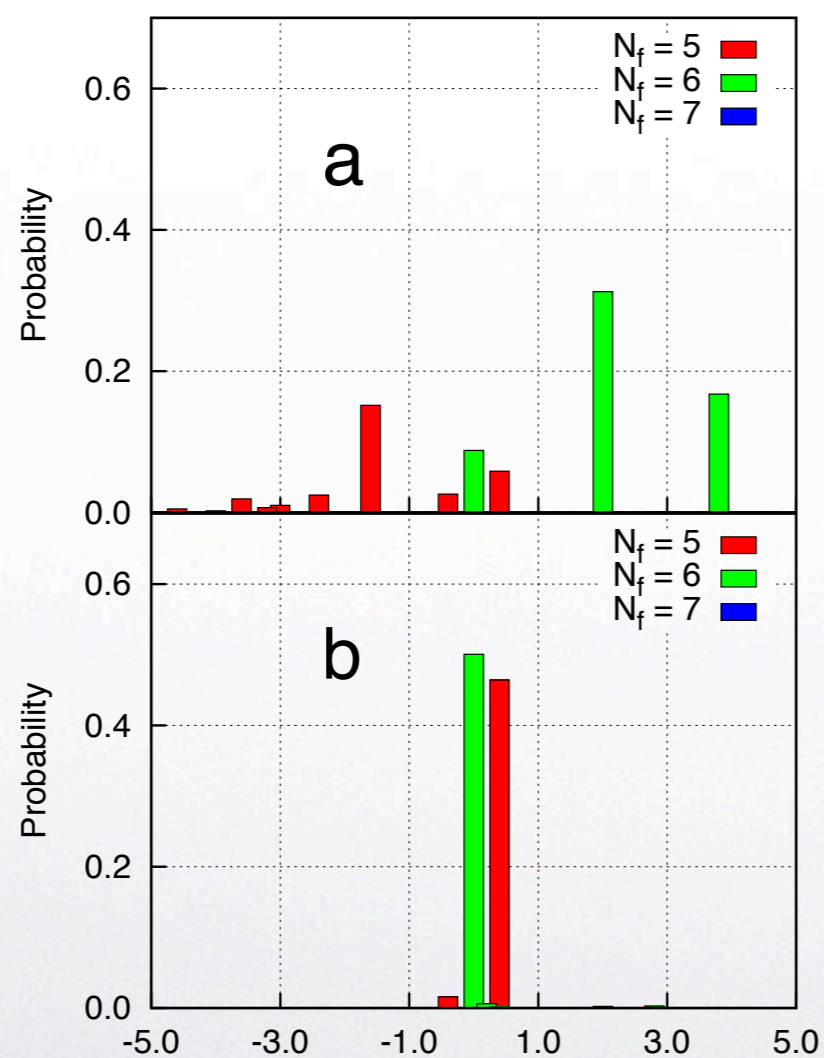


# Band structure obtained by LDA and LDA+G





# The probability of atomic eigenstates for SmB6

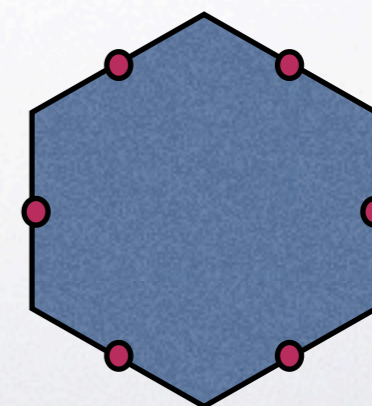
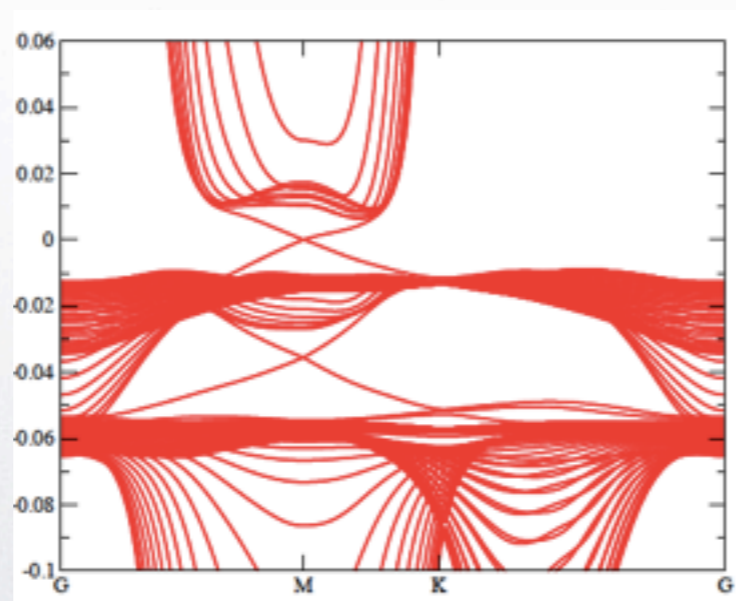
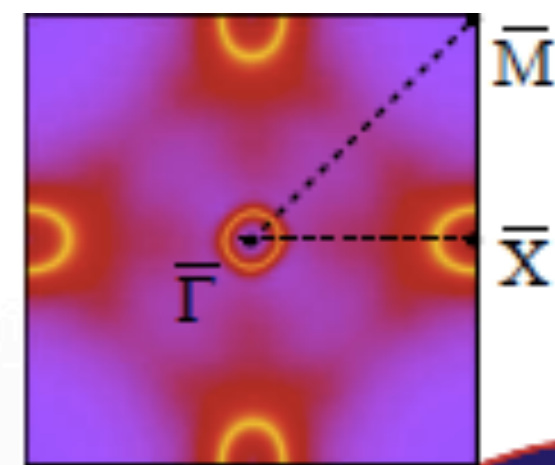
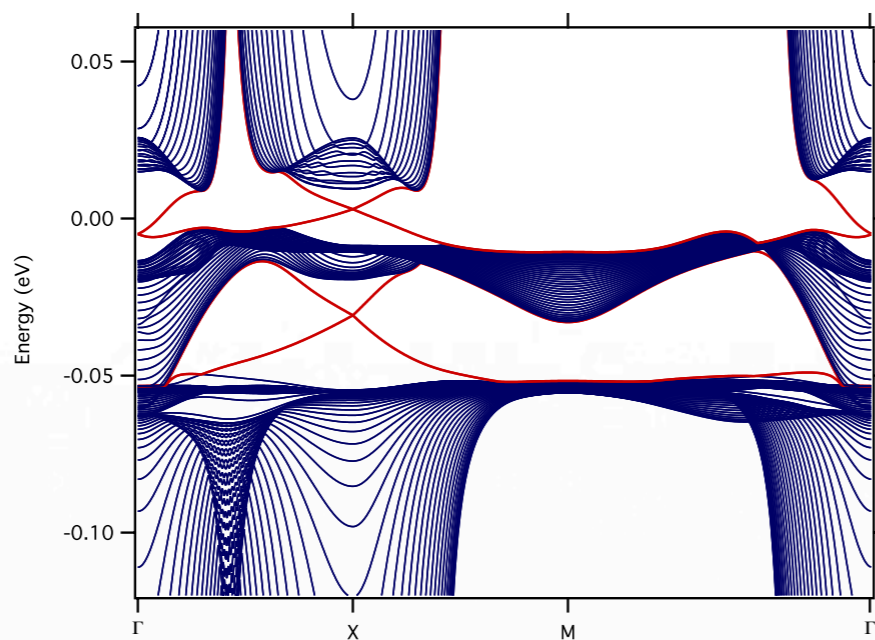


LDA

LDA+G

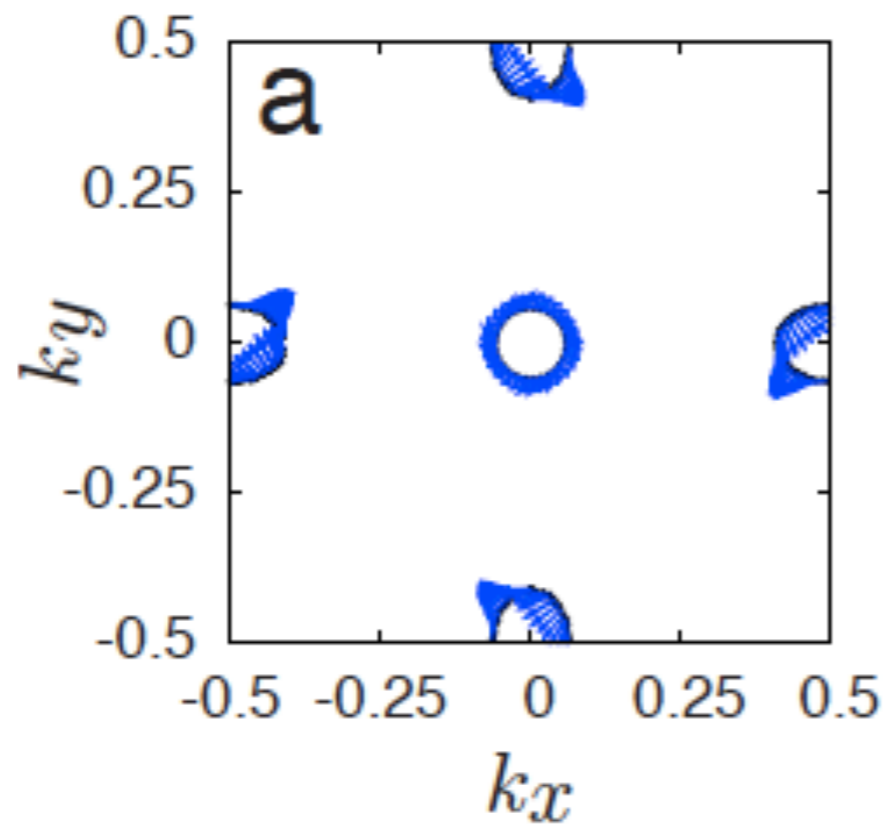


# The unique surface states of SmB6 on (001) surface

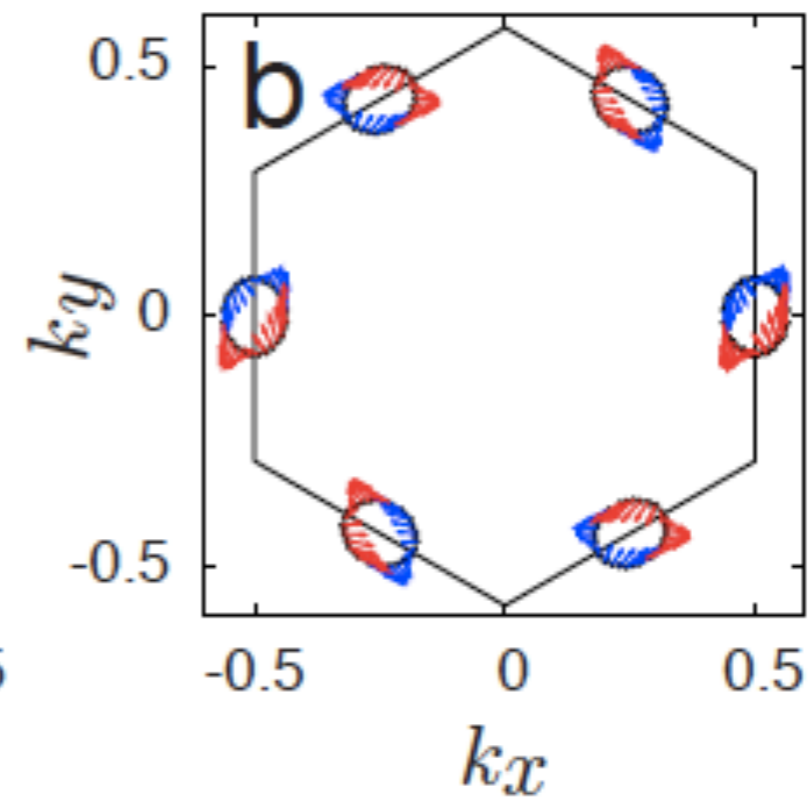




## Spin texture for surface states



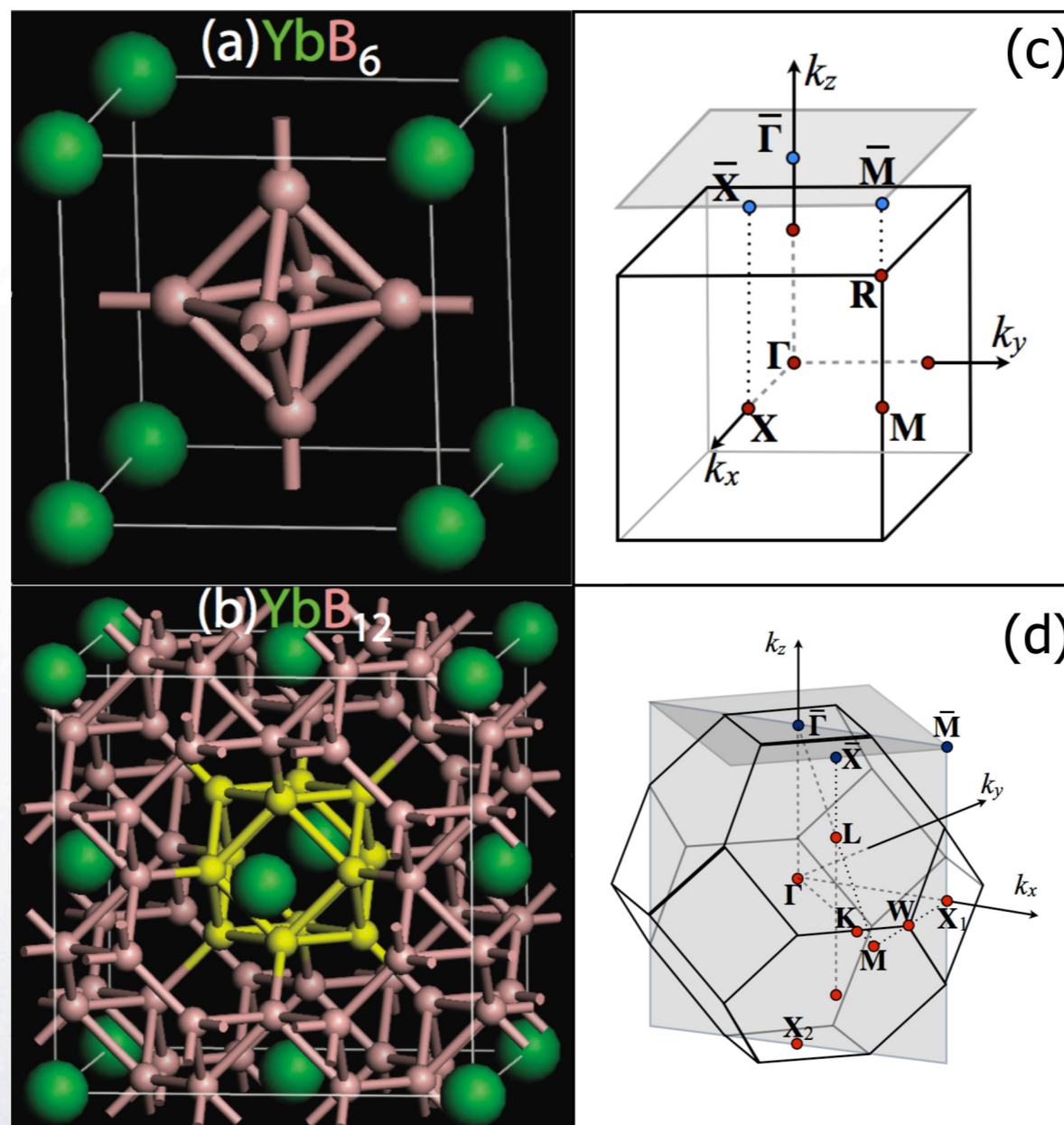
001 surface



111 surface



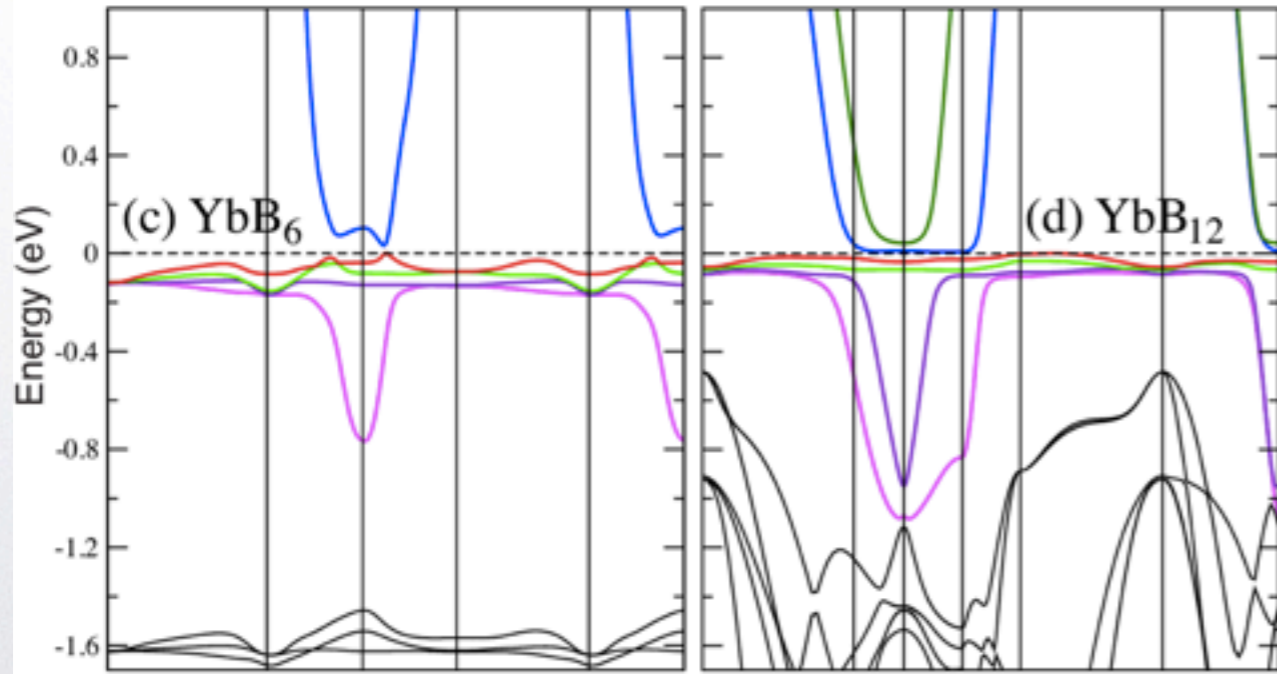
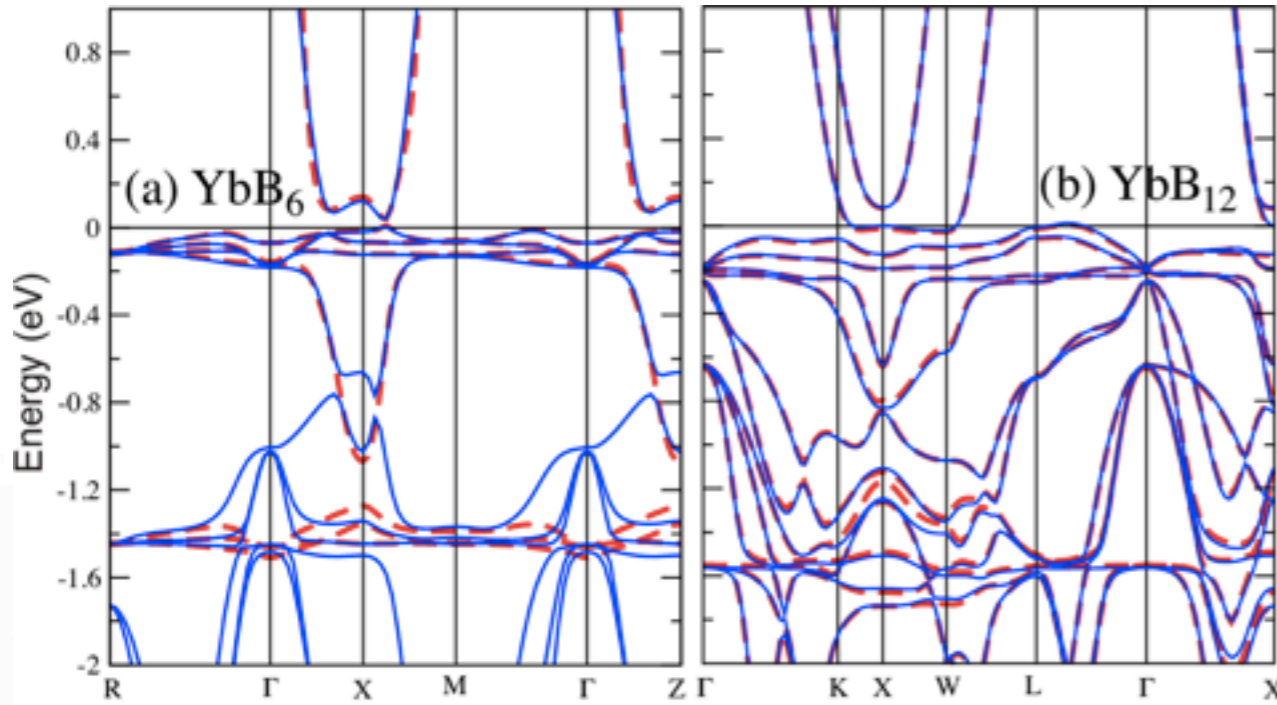
# Structure of YbB<sub>6</sub> and YbB<sub>12</sub>



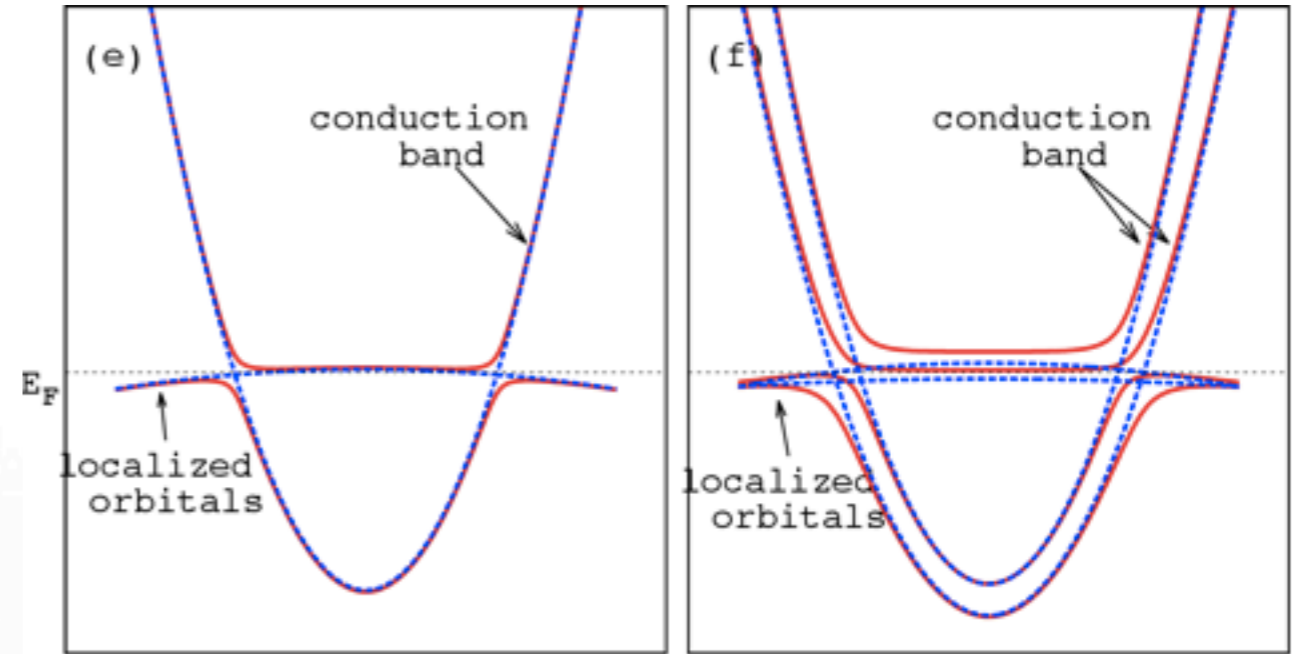
# Renormalized Band structure



## LDA



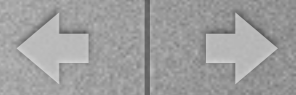
## LDA+Gutzwiller



	$\Gamma$	$3X$	$3M$	$R$	$n_f$	$z$
YbB <sub>6</sub>	+	-	+	+	13.80(13.58)	0.87
YbB <sub>12</sub>	+	+	+	+	13.11(13.31)	0.28



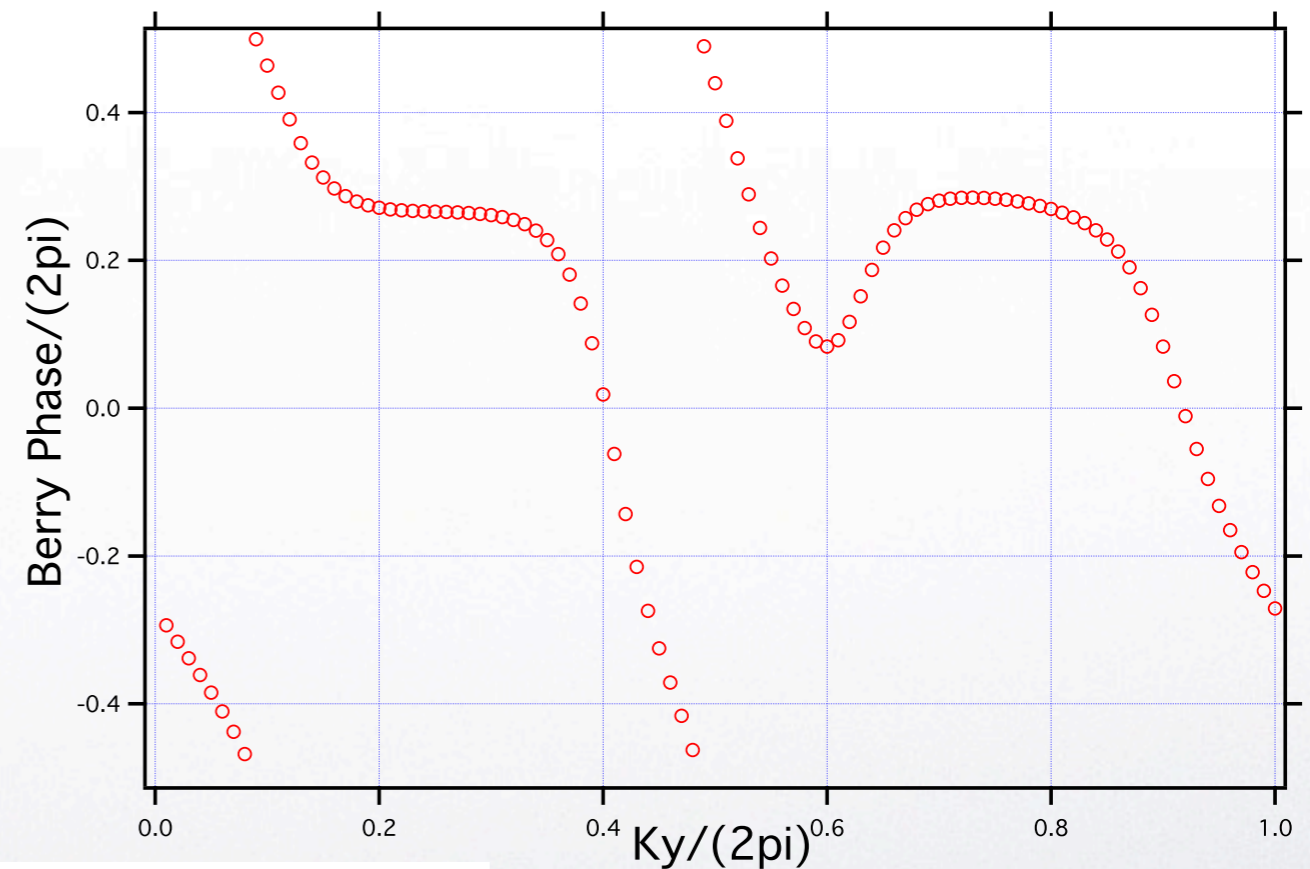
# Determine the Mirror Chern number



By the eigenvalue of the  $C_n$  rotational at high symmetry point

	C4	C2	M
4f7/2 $j_z=1/2$	$\exp(i\pi/4)$	$\exp(i\pi/2)$	-i
5d5/2 $j_z=-1/2$	$\exp(-i\pi/4)$	$\exp(-i\pi/2)$	-i
4f7/2 $j_z=-3/2$	$\exp(-i3\pi/4)$	$\exp(-i3\pi/2)$	-i
5d5/2 $j_z=3/2$	$\exp(i3\pi/4)$	$\exp(i3\pi/2)$	-i

By Wilson Loop method



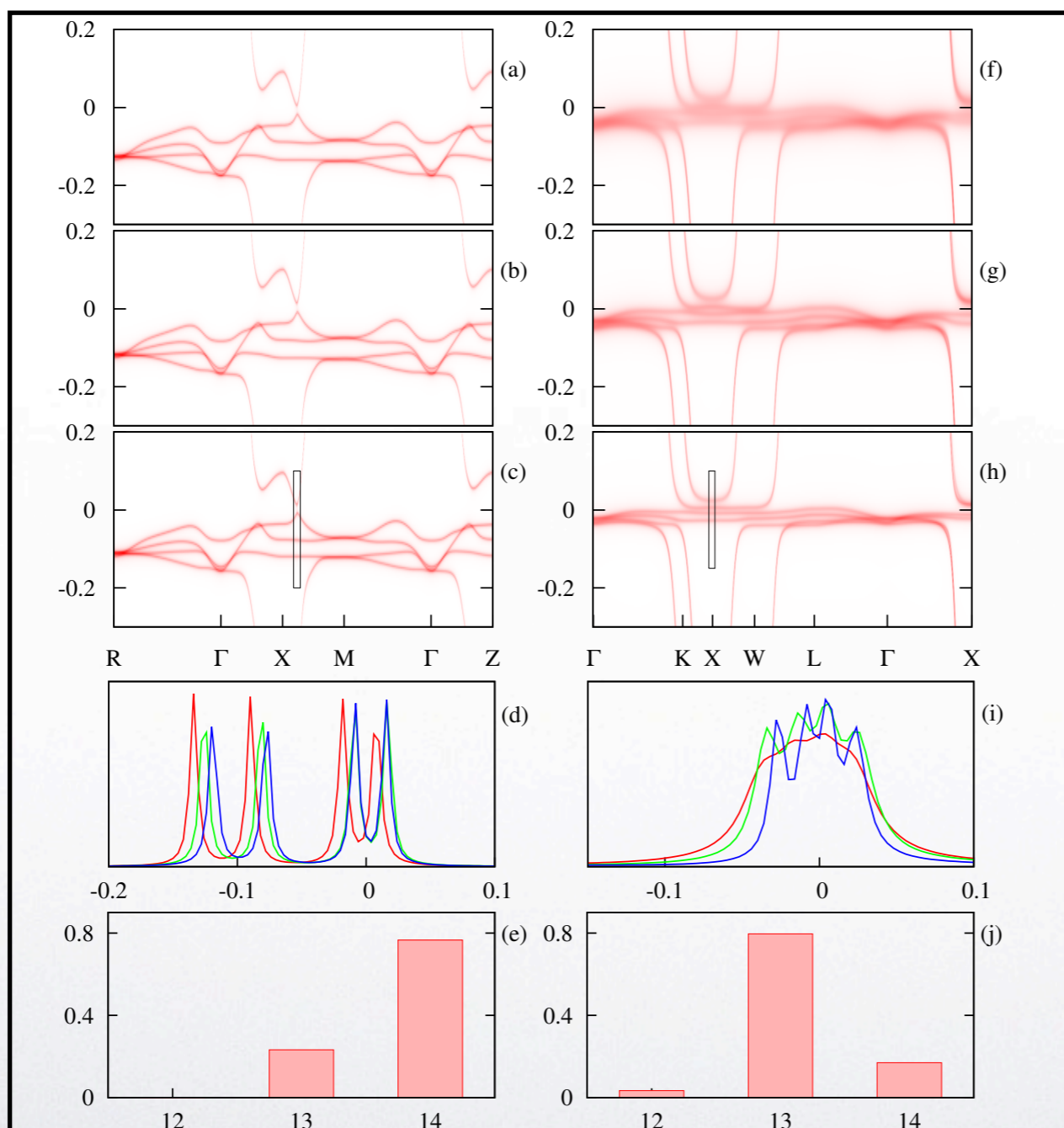
$$i^C = -1$$

$$\exp(i2\pi C/4) = \prod_{n \in \text{occ}} (-1)^F \xi_n(\Gamma) \xi_n(M) \zeta_n(X), \quad (5)$$

C. Fang et al, PRL 108, 266802 (2012); R. Yu, et al, Phys. Rev. B **84**, 075119 (2011)



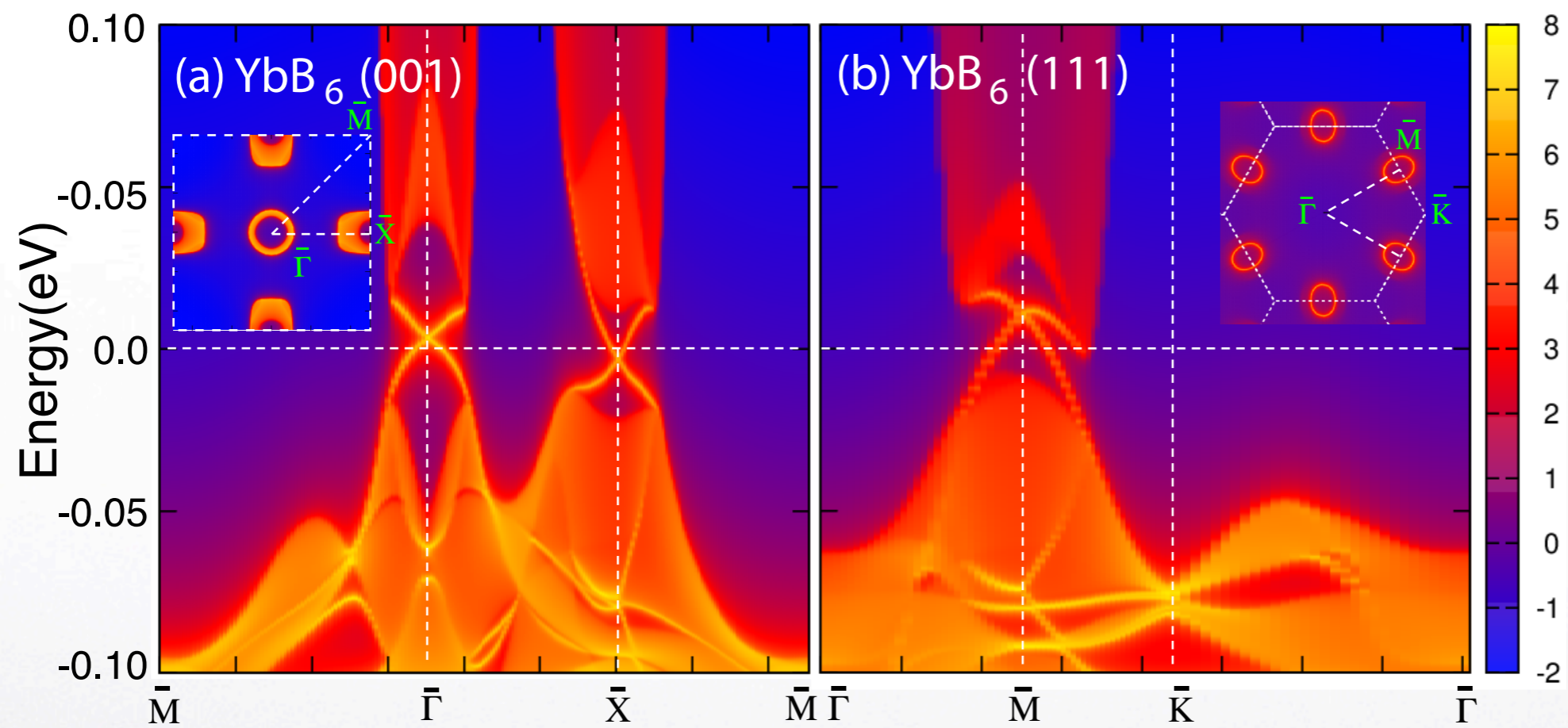
# Temperature dependance of the spectral function





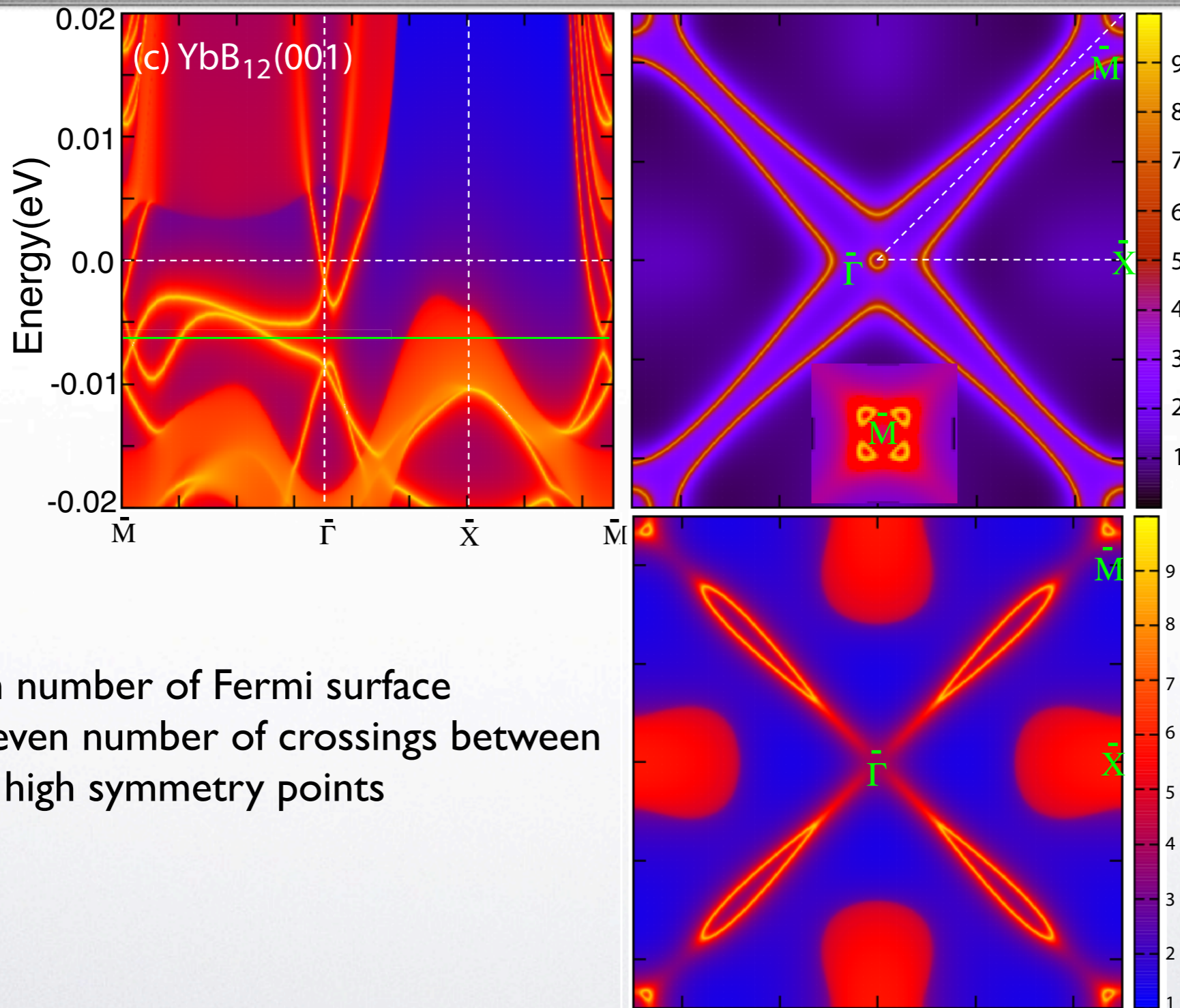
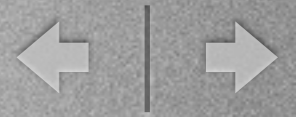


# Surface states for YbB<sub>6</sub>





# Surface states for YbB<sub>12</sub>





# Conclusions

- Mix valence TI SmB<sub>6</sub>, and YbB<sub>6</sub>
- Topological Crystalline Kondo Insulator with mirror Chern number 2 : YbBI<sub>2</sub>
- The topological phases survive from the Strong e-e interaction among f electrons
- More similar materials!